

Space Charge Resonances in High Intensity Drivers

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I. Hofmann, A. Fedotov (BNL), G. Franchetti

GSI Darmstadt

Outline

- Coherent Space Charge Tune Shift
 - Analytical Approach
 - Coupled Mode Coefficient
- Lattice Driven Resonances (2D)
 - Half-Integer Resonance (SNS)
- Space Charge Driven Resonances

Require systematic approach

1. For the GSI project of a large synchrotron SIS100 as driver for radioactive ions and antiprotons we need to keep loss $<1\%$ (~ 1 sec) - main problem seems wall desorption (10^4 ions per incident heavy ion)
1. 3 kinds of space charge effects:
 - tune spread/shift -> „frozen-in“ space charge for up to $\sim 10^5$ turns (talk by G. Franchetti)
 - sc as source of nonlinearity
 - coherent dynamical effect -> self-consistent space charge calculation for $\sim 10^3$ turns - compare with theory (this talk); applies to coasting beam or ideal barrier bucket

Hofmann, Beckert, 1984

Machida, 1991 ->

Baartman, 1998

Venturini, Gluckstern, 2000

Burov, 2001

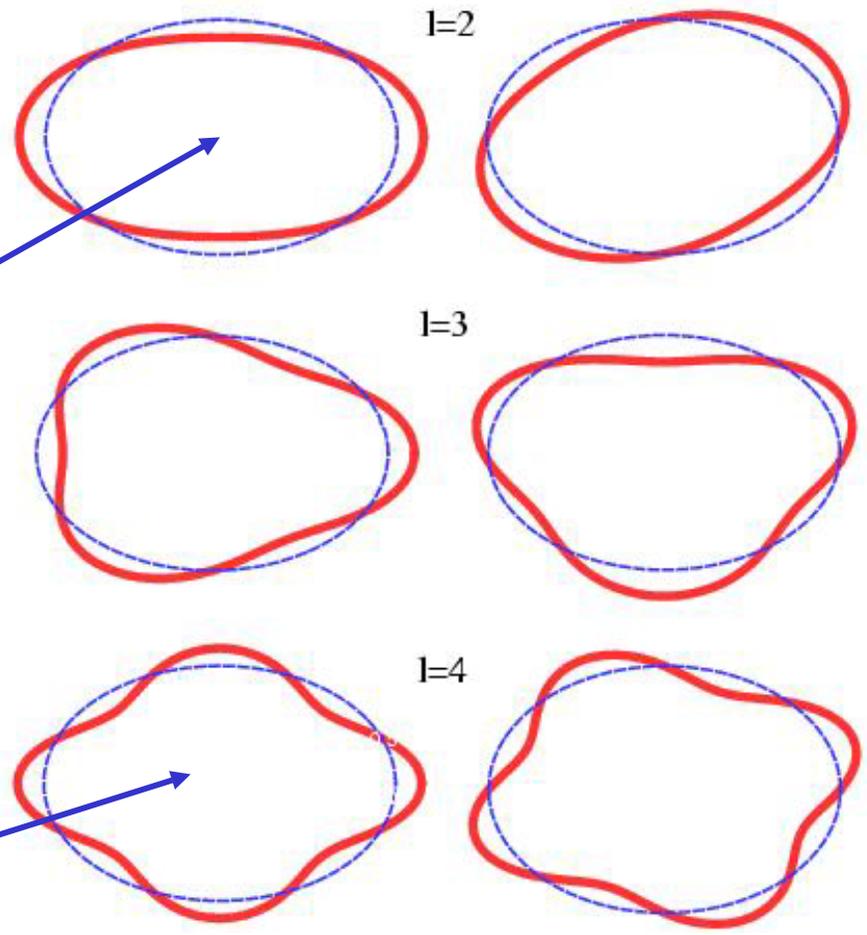
Fedotov, Hofmann, 2002

Transverse Space Charge Modes

Eigen modes of transverse density oscillations
Coherent motion induces additional shift
-> needs selfconsistent (kinetic) calculation

normal
even modes

skewed
odd modes



2 envelope modes
(breathing and quadrupolar)

Normal octupole or
Space charge of WB beam

Derived dispersion relation from Vlasov-Poisson equation for $l = 2, 3, 4$

Hofmann., *Phys. Rev. E* **57** (1998)

The coherent frequency $\sigma = \omega/v_x$ must be solved in 3-dimensional parameter space:

$$\alpha = v_y/v_x, \eta = a/b, \sigma_p = \omega_p/v_x \text{ with } \omega_p^2 = q^2 N / (\epsilon_0 \pi m \gamma^3 ab)$$

Example:
 $l=3$, even

Isotropic:

$$D_{3,e} = 8 + \sigma_p^2 \frac{12}{9 - \sigma^2} - \sigma_p^4 \frac{4\sigma^2(3 - \sigma^2)}{(9 - \sigma^2)^2(1 - \sigma^2)^2}$$

Anisotropic:

$$\begin{aligned} D_{3,e} \equiv & (1 + \eta)^3 + \\ & \frac{\sigma_p^2}{8} \left[\frac{1 - 5\eta}{1 - \sigma^2} + \frac{9 + 27\eta + 24\eta^2}{9 - \sigma^2} \right. \\ & \left. + \frac{(1 - 2\alpha)(1 - 2\eta^2/\alpha)(3 + \eta)}{(1 - 2\alpha)^2 - \sigma^2} + \frac{(1 + 2\alpha)(1 + 2\eta^2/\alpha)(3 + \eta)}{(1 + 2\alpha)^2 - \sigma^2} \right] + \\ & \frac{\sigma_p^4}{8} \left[\frac{-1}{(1 - \sigma^2)^2} + \frac{3}{(1 - \sigma^2)(9 - \sigma^2)} \right. \\ & \left. + \frac{3(1 - 2\alpha)}{(9 - \sigma^2)((1 - 2\alpha)^2 - \sigma^2)} + \frac{3(1 + 2\alpha)}{(9 - \sigma^2)((1 + 2\alpha)^2 - \sigma^2)} \right] \\ & = 0 \end{aligned}$$

Im $\sigma > 0$:
many resonant
instabilities

Recently resolved confusion about KV and non-KV modes & instabilities

„Gluckstern modes“ in symmetric beams known as KV artifact

- **Anisotropic**: we claimed some modes are „real“ (non-oscillatory) – not unambiguous
- Others emphasise highest frequency branches („fluid modes“ by Lund&Davidson“)

Key is „**coherent coupled mode coefficient**“, generalizing Baartman's definition for symmetric beams to unsymmetric/anisotropic beams:

$$\omega = m v_z + l v_x + \Delta\omega = n$$

$$C_{mlk} = (m v_{0z} + l v_{0x} - \omega) / (m \Delta v_z + l \Delta v_x)$$

gives full coherent shift versus „incoherent coupled mode shift“

$$C_{202} \text{ (envelope)} = 0.75 \rightarrow 33\% \text{ „intensity gain“}$$

new finding: $C > 1$ (shift past incoherent) are KV-specific „negative energy modes“ leading to „Gluckstern“ instabilities - need to be discarded

$C < 1$ are positive energy modes of realistic beams \rightarrow anisotropy instabilities for $C = 1$
 \rightarrow most low-frequency branches also realistic (kinetic – not fluid)

Example:

$$v_{0z} = 4.115 / v_{0x} = 4.35$$

$$\epsilon_z / \epsilon_x = 2$$

4th order modes

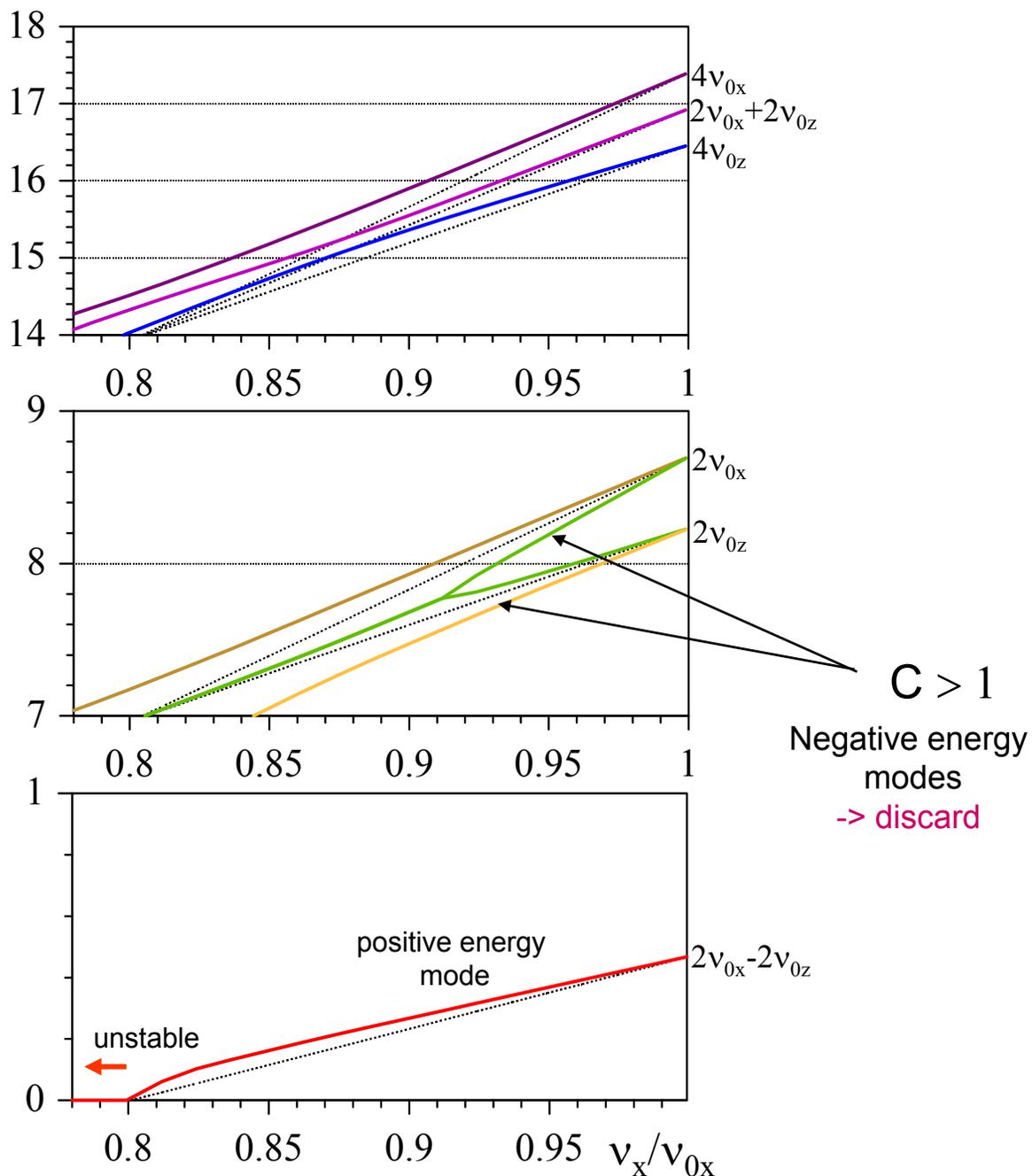
Examples for highest frequency branches:

2nd order: $C_{20} = 0.65$

3rd order: $C_{30} = 0.77$

4th order: $C_{40} = 0.88$

trend: approaching 1

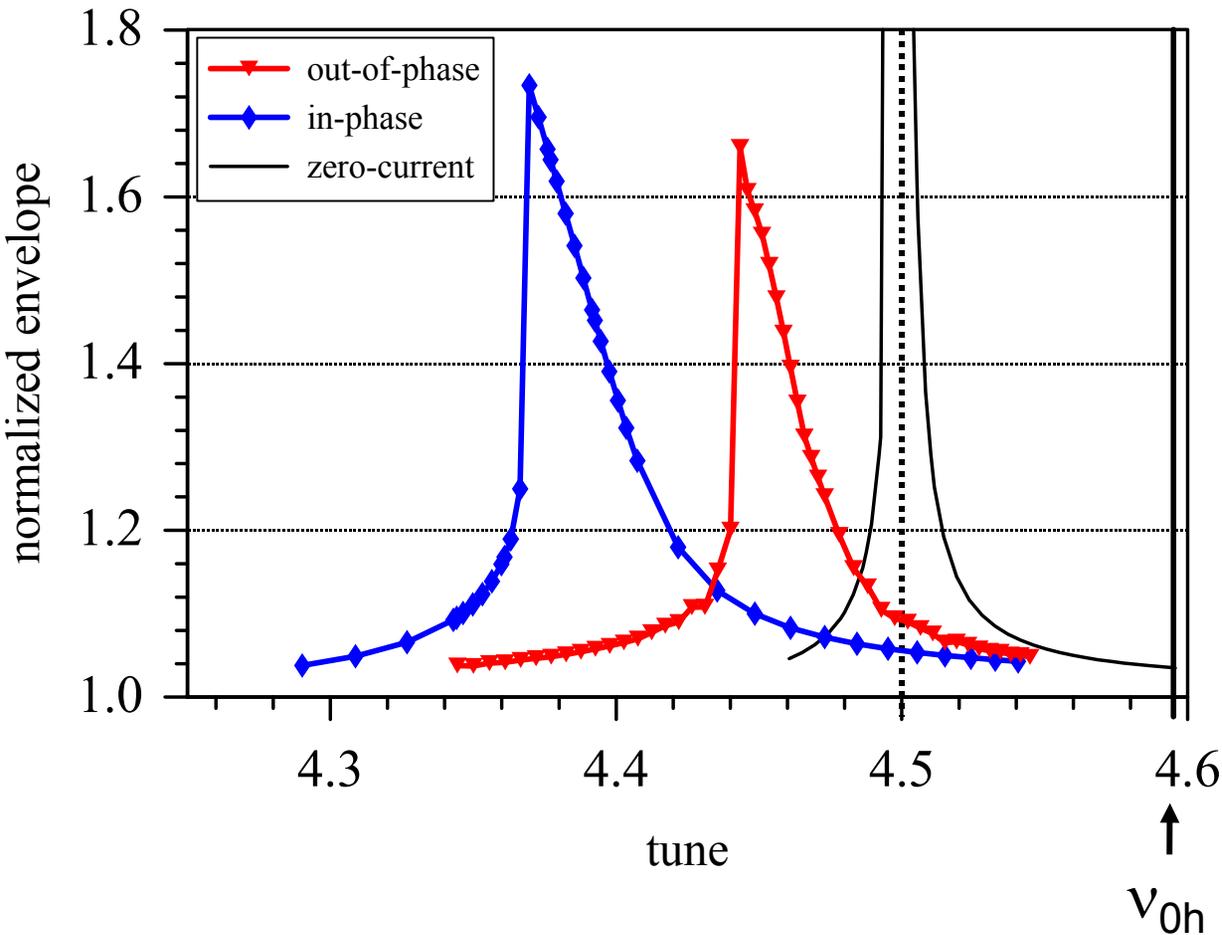


Applied to SNS ring: essentials from envelope equation

2nd order even mode resonance: $2\nu_h + \Delta\omega = 9$

driven by gradient error near (fictitious) working point 4.5

Fedotov, Hofmann, PRSTAB 2002



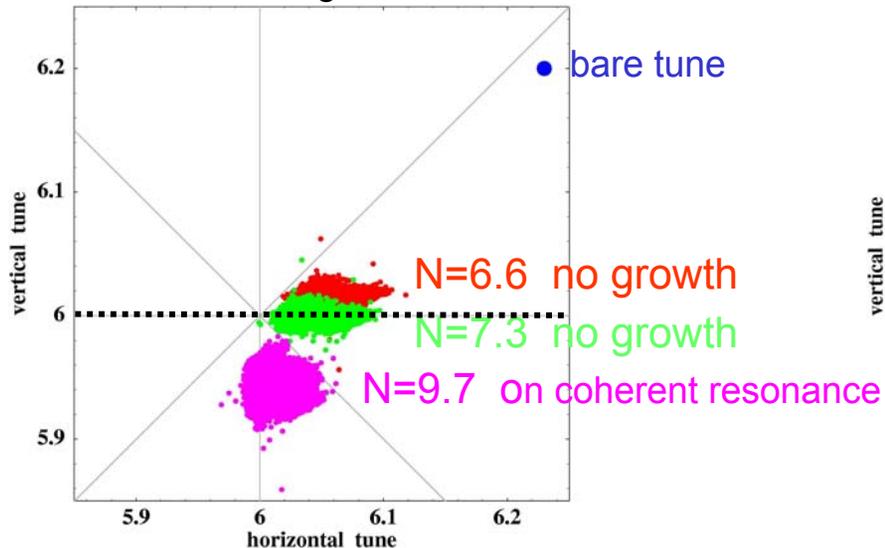
- **coherent shift** of resonance
- **different** for breathing and quadrupolar mode
- envelope response limited due to **de-tuning**

next: applied ORBIT code to SNS lattice

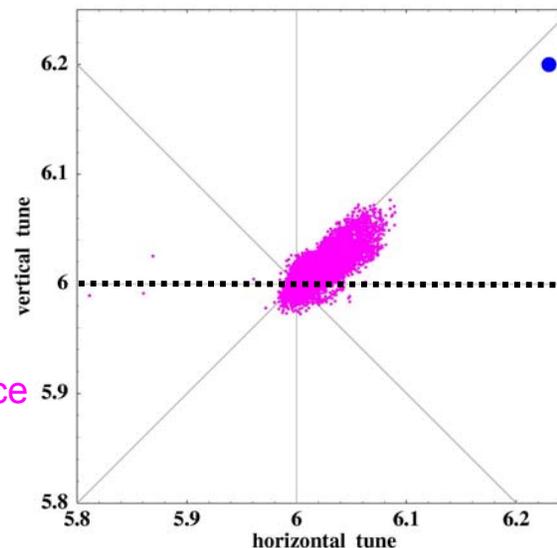
Found clear intensity gain due to the coherent shift of the 2nd order (systematic) resonance $2\nu_v + \Delta\omega = 12$ (4 super periods)

intensity gain basically maintained for injection into bucket (< 1 synchrotron period)

single particle footprints (1st turn)
for initial coasting KV beam

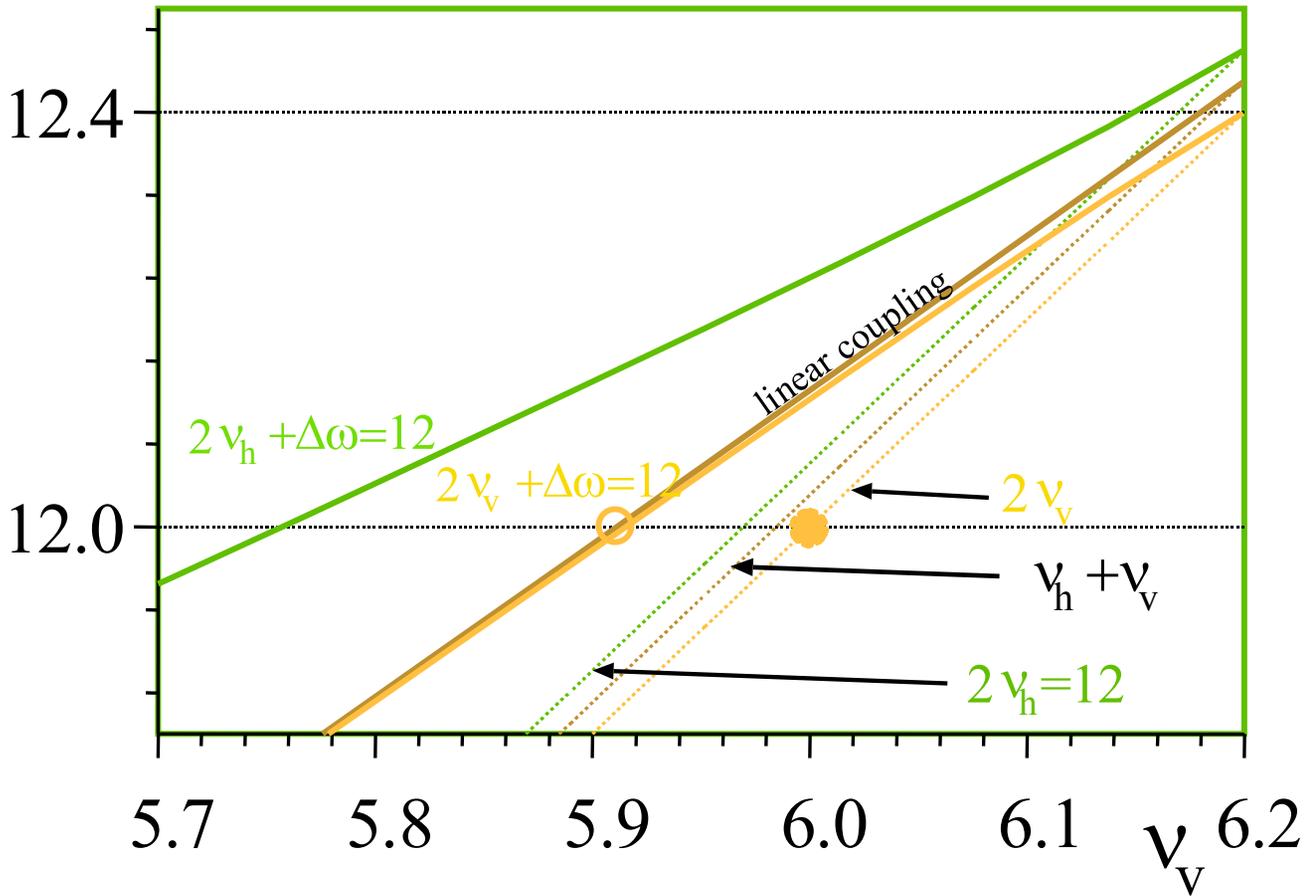


100 turns



Consistent with theoretical shift

$$2\nu_v + \Delta\omega = 12$$



$\nu_0 = 6.23/6.20$; emittance ratio=1

Check 4th order:

fourth order resonances much
closer to $2\nu_v=12$:

$$4\nu_v + \Delta\omega = 24$$

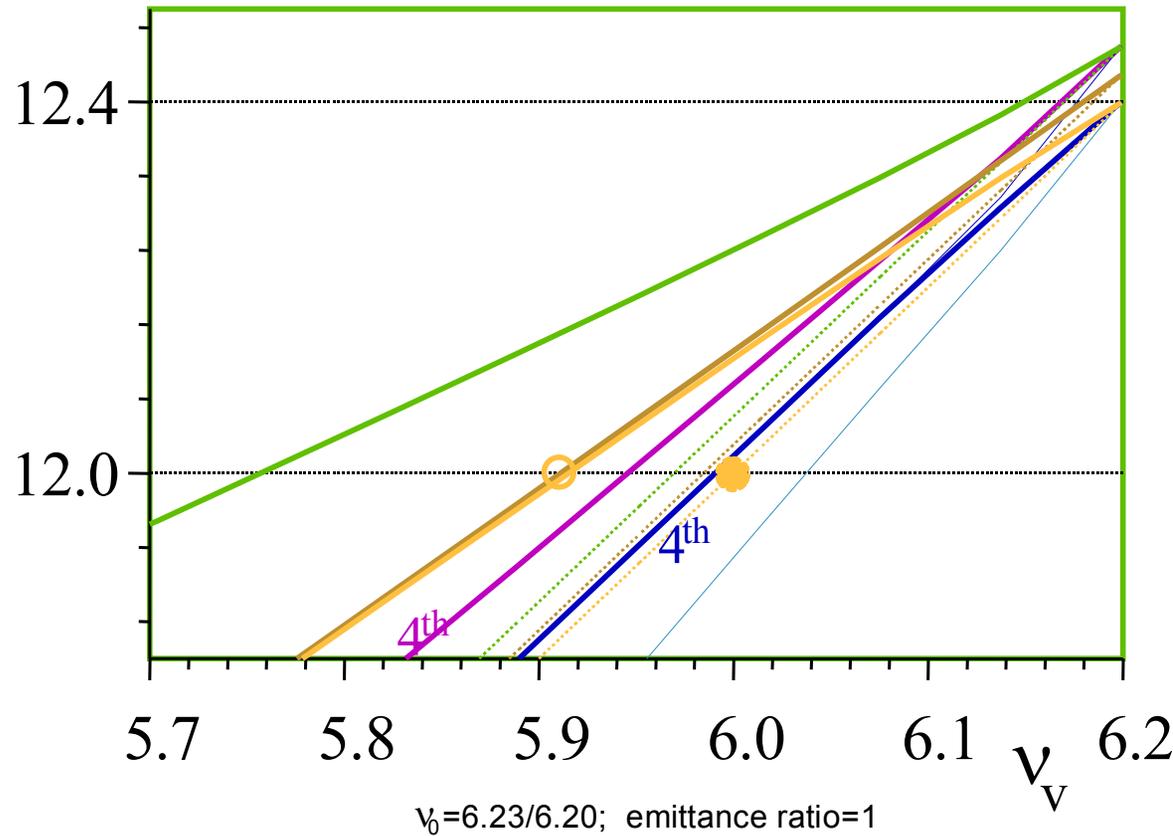
and:

$$2\nu_v + \Delta\omega = 12$$

does this take away „coherent
advantage“?

first tests turning on fringe
field octupoles in SNS: no

but: higher order might show
up $\gg 1\text{ms}$
need careful study!



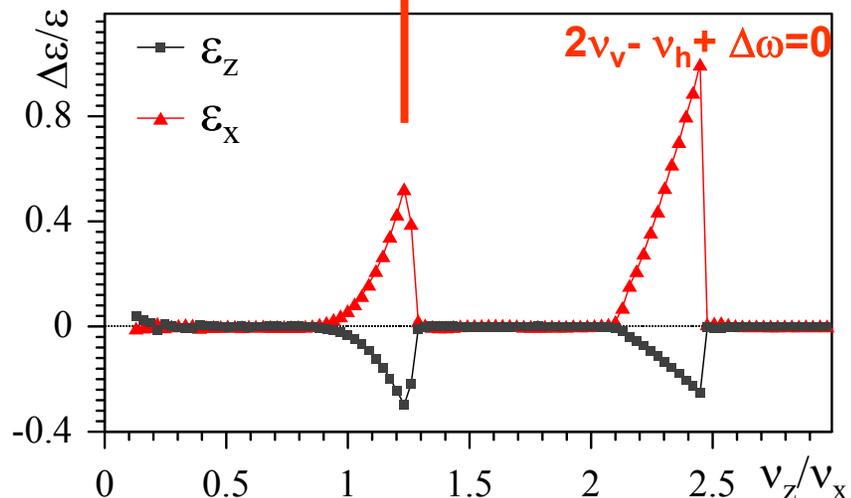
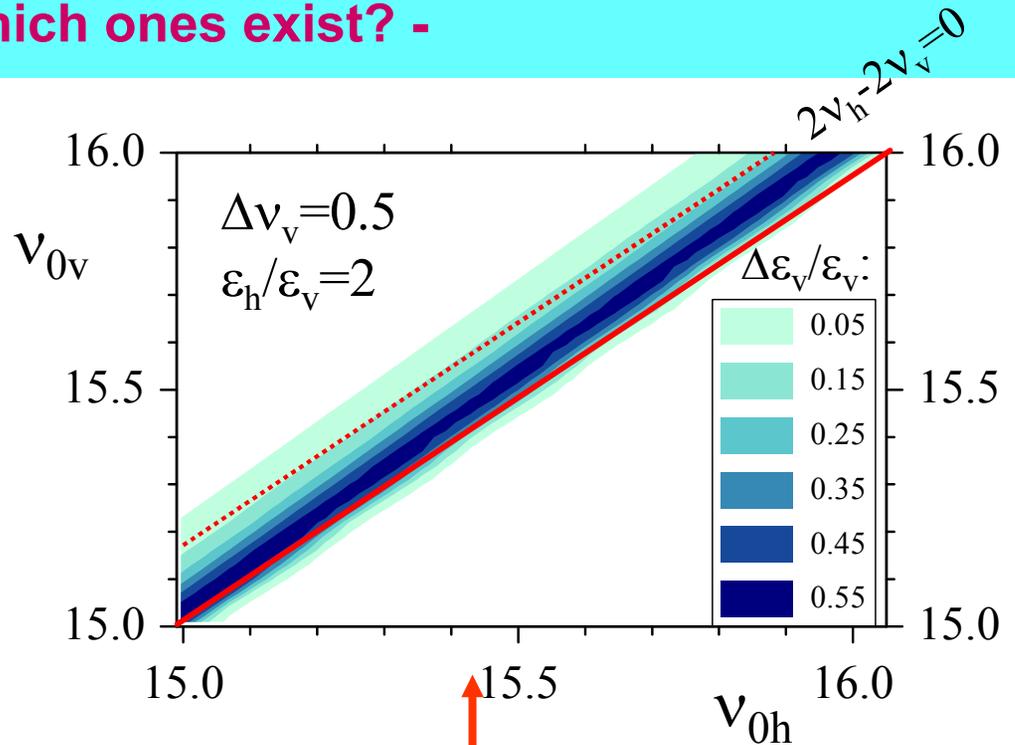
Higher order space charge driven resonances in linear lattice - which ones exist? -

„Montague“ resonance:
 $2v_v - 2v_h + \Delta\omega = 0$

expect also:
 $2v_v + 2v_h + \Delta\omega = n \times N$
studied by Machida et al.

data taken from scan over stopband
 using 2D PIC simulations for initial waterbag:

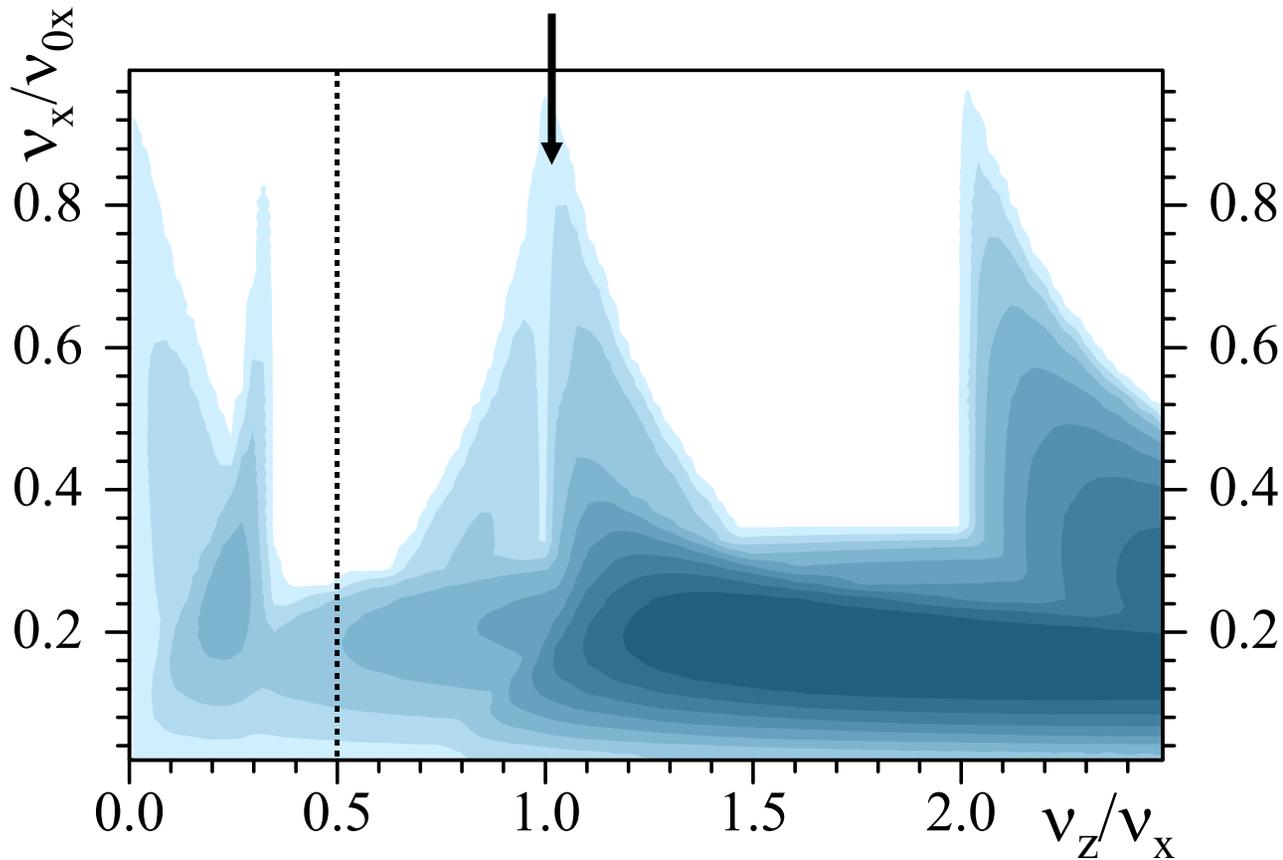
But: also find resonant growth
 (exponential) in third order due to
 space charge - we claim not only
 even sc-resonances matter!
 Found $2v_h + v_v + \Delta\omega_3 = 12$ in pure
 linear lattice!



In Linacs same stopband responsible for emittance exchange

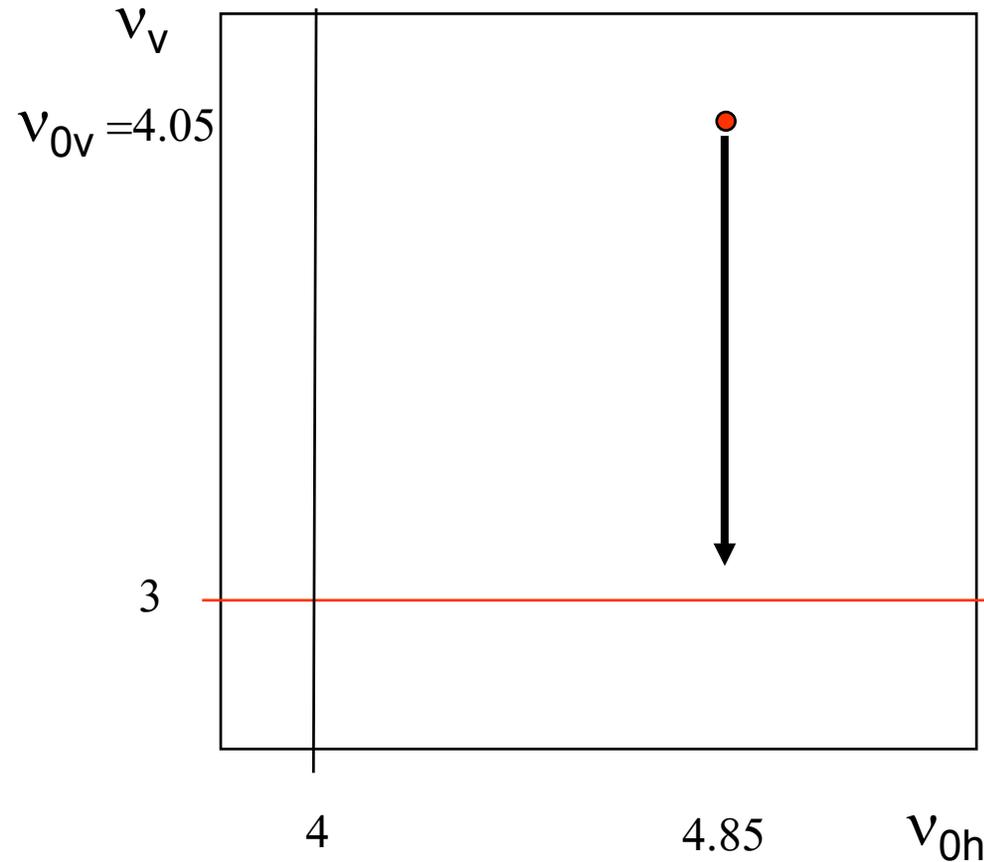
„Montague resonance“

Montague, 1968
Hofmann, Boine-Frankenheim, PRL2001)
Sakai et al. PAC01

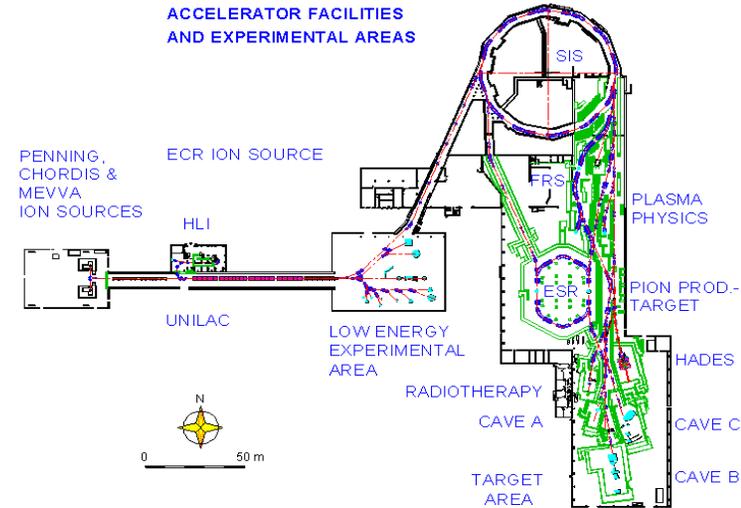


2d self-consistent PIC simulations for SIS18 lattice (FODO with 12 super-periods)

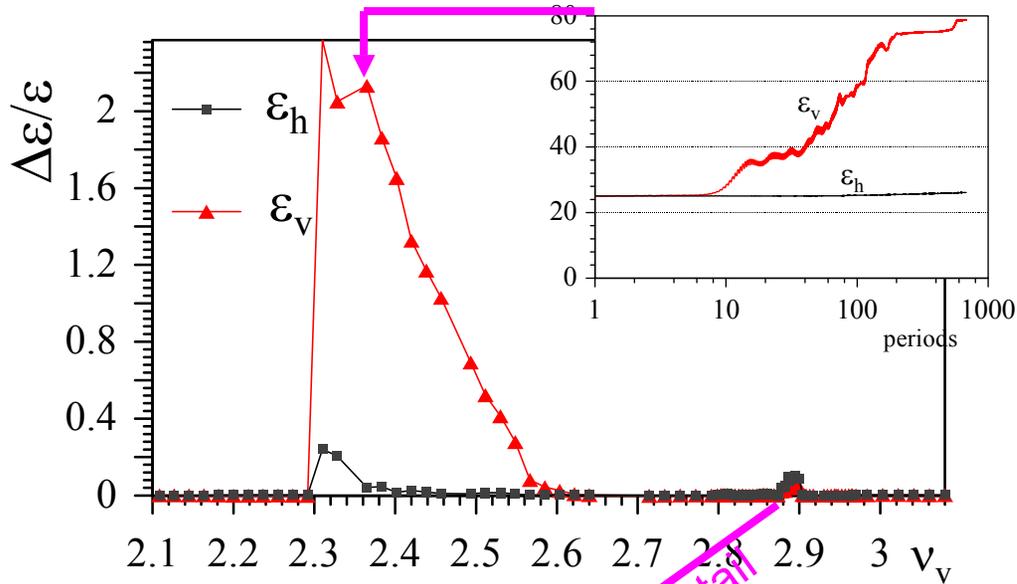
using initial waterbag in 2D coasting beam/ 50.000 particles PIC



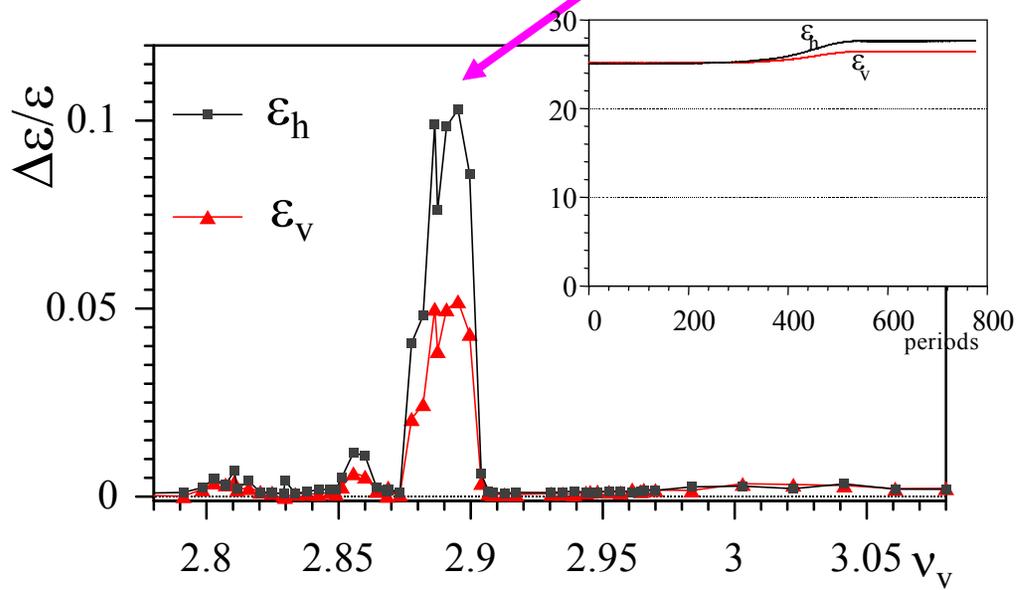
„ $4v_v = 12$ “



Found modes growing from noise = exponential instabilities



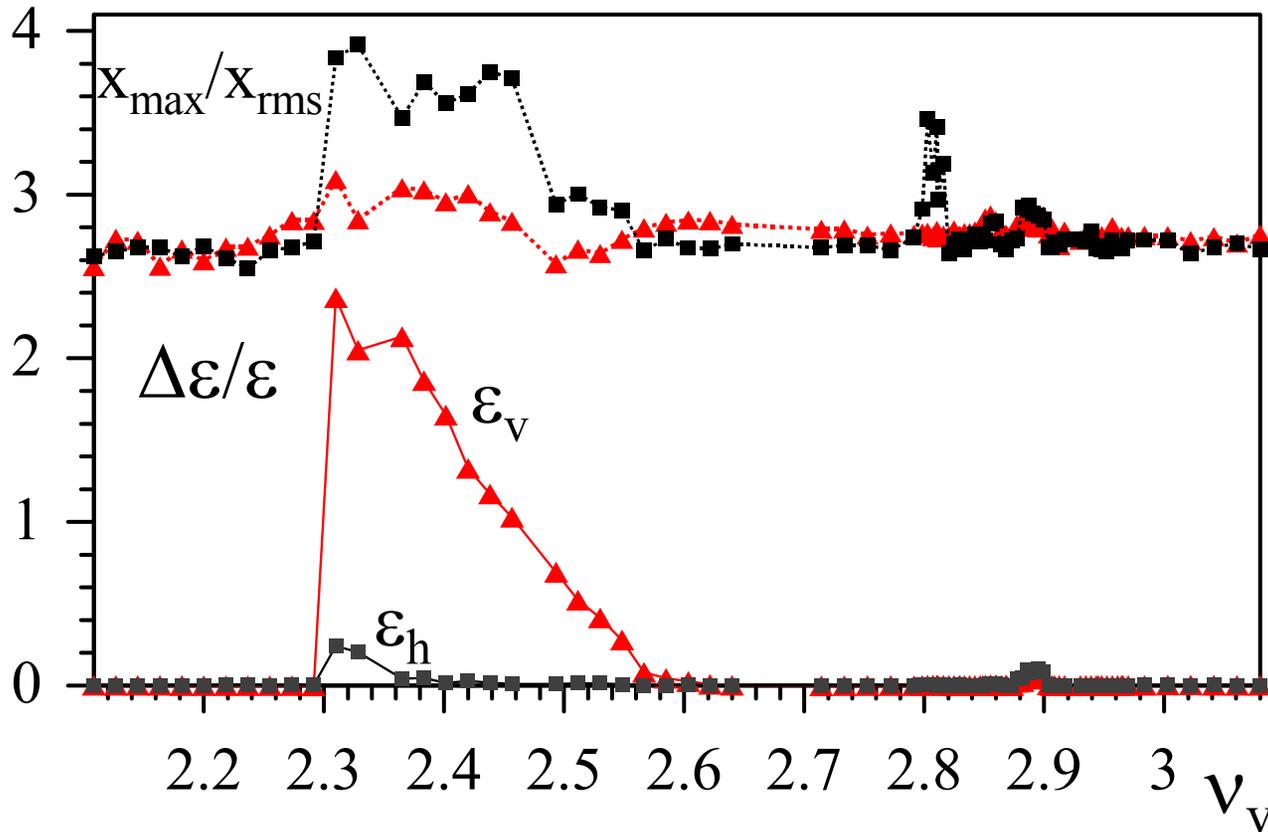
$2\nu_V + \Delta\omega_2 = 12/2$
 envelope instability with $\sigma_0 > 90^\circ$
 completely absent without space charge
Hofmann, Laslett, Smith, Haber, 1983



fourth order:
 $4\nu_V + \Delta\omega_4 = 12$
 $2\nu_V + \Delta\omega_4 = 12/2$
 $C \sim 0.88$ is in good
 agreement with theory

- **coherent shift** of resonance
- **emittance growth limited** due to **sc de-tuning**

Maximum radius



- Rms emittance growth limited due to detuning
- Halo larger where less rms emittance growth

Conclusions

1. Coherent resonance crossing & shift in second order
 - favourable coherent tune shift for \sim coasting beam (SNS)
2. For higher order modes found selection rule to discard „KV-modes“ ($C_{\text{mlk}} > 1$)
 - coherent shift smaller in fourth order
3. Also found systematic sc driven instabilities $m\nu_h + l\nu_v + \Delta\omega = N/2$ in all orders
4. Pronounced self-limiting (de-tuning) effect of space charge in nonlinear ($> 2^{\text{nd}}$ order) resonance
 - suggests quite small resonance loss in 2D coasting beam or ideal barrier bucket -> significant resonance advantage of barrier bucket
5. Require „sufficiently fast“ synchrotron motion to explain observed losses in synchrotrons